

Nonlinear surface magneto-optics of Fe/Cu(001) from first principles: Influences of the Cu substrate

Torsten Andersen

Condensed Matter Theory Group, Ångström laboratory
Uppsala University, P.O.Box 534, SE-75121 Uppsala, Sweden

Max-Planck-Institut für Mikrostrukturphysik
Weinberg 2, D-06120 Halle (Saale), Germany

thor@mpi-halle.mpg.de

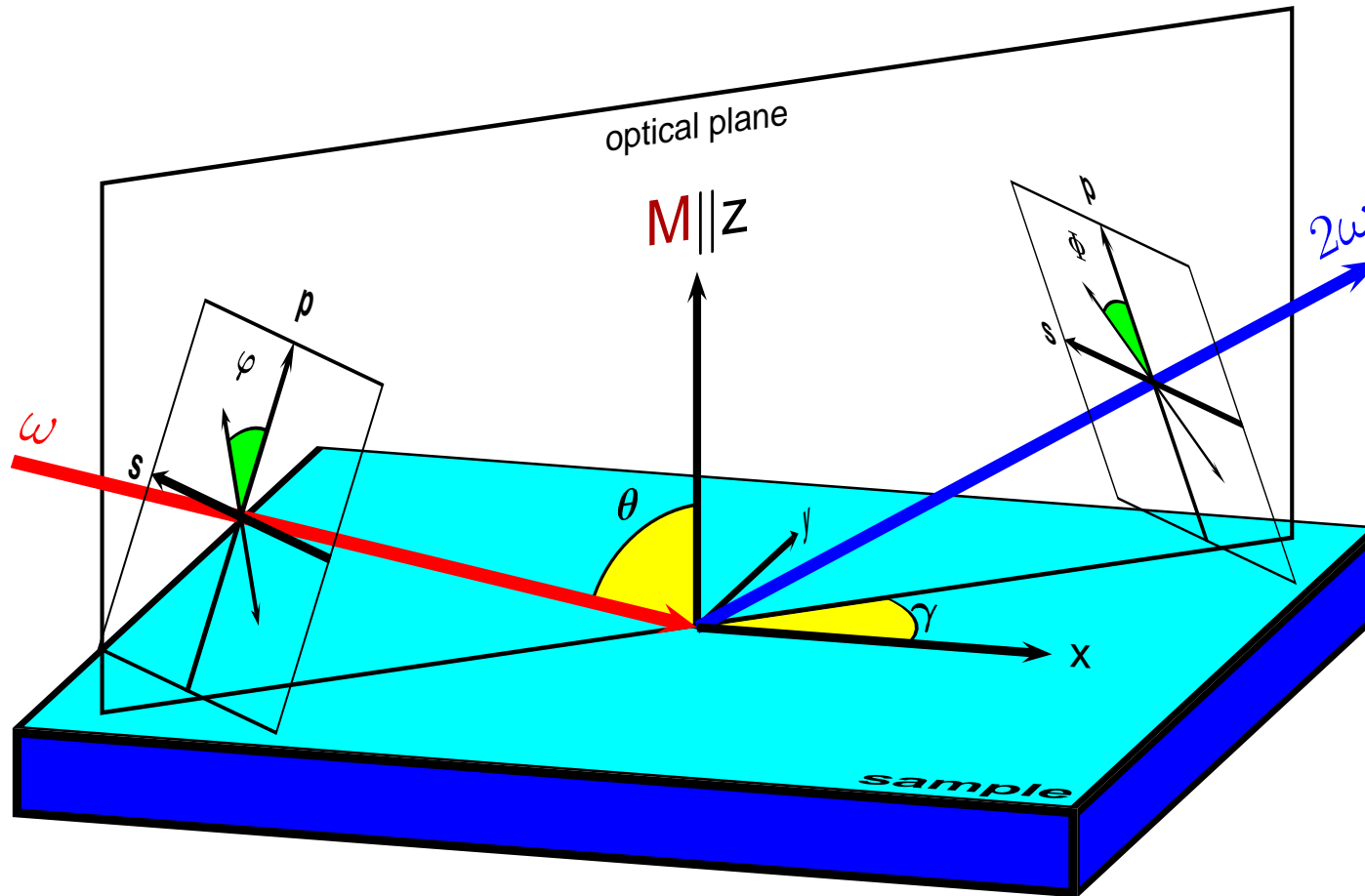
Acknowledgements

- Vector optimization: Rechenzentrum Garching
- Financial support: DFG and EU TMR "NOMOKE" and "DYNASPIN"

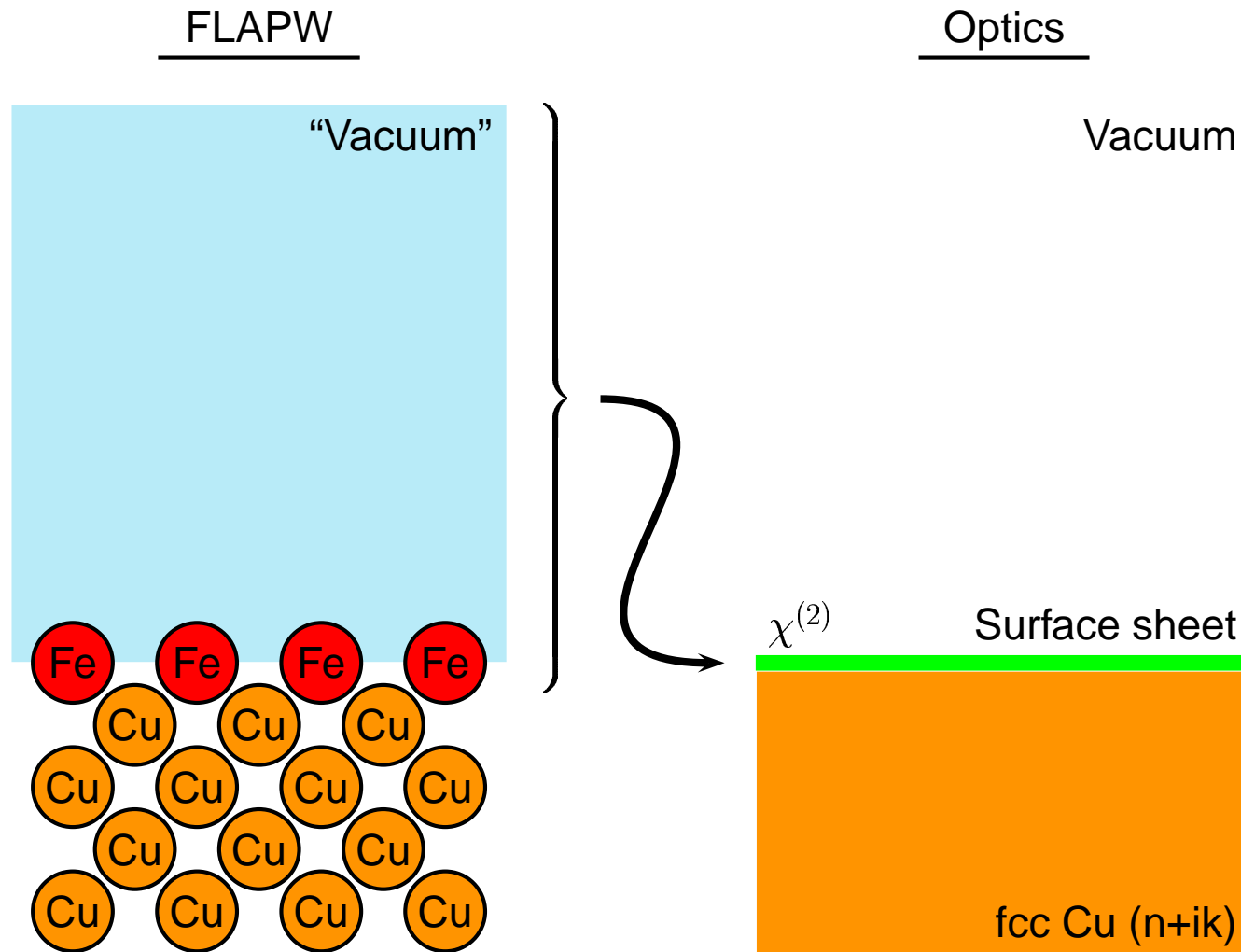
Outline

- Reflection geometry for nonlinear optics
- Numerical calculations: ▶ method ▶ formulas
- Results: ▶ band structure ▶ layer-resolved magnetic moments ▶ screened nonlinear susceptibility ▶ nonlinear optical intensities ▶ rotating the polarizer ▶ nonlinear Kerr rotation
- Conclusions and outlook

Reflection geometry for SHG on magnetic surface



Nonlinear magneto-optics – method



Nonlinear optical field

$$E(2\omega; \theta, \Phi, \varphi, \gamma) = 2i\delta z \frac{\omega}{c} |E_0(\omega)|^2 \chi_{ijk}^{(2)}(2\mathbf{q}, 2\omega), \mathbf{M} \parallel \mathbf{z}$$

$$\times \left\{ A_s \sin \Phi \begin{pmatrix} -\sin \gamma \\ \cos \gamma \\ 0 \end{pmatrix} + A_p \cos \Phi \begin{pmatrix} F_c \cos \gamma \\ F_c \sin \gamma \\ N^2 F_s \end{pmatrix} \right\} \cdot \begin{pmatrix} 0 & 0 & 0 & | & xyz^- & xxz^+ & 0 \\ 0 & 0 & 0 & | & xxz^+ & -xyz^- & 0 \\ zxx^+ & zxx^+ & zzz^+ & | & 0 & 0 & 0 \end{pmatrix}$$

$$\cdot \begin{pmatrix} [t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma]^2 \\ [t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma]^2 \\ t_p^2 f_s^2 \cos^2 \varphi \\ 2[t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma] t_p f_s \cos \varphi \\ 2[t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma] t_p f_s \cos \varphi \\ 2[t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma][t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma] \end{pmatrix}$$

Screened nonlinear optical susceptibility

$$\chi_{ijk}^{(2)}(2\mathbf{q}, 2\omega) = e^3 \sum_{\mathbf{k}, l, l''} \left\{ \frac{\langle \mathbf{k} + 2\mathbf{q}, l'' | \mathbf{r}_i | \mathbf{k}, l \rangle \langle \mathbf{k}, l | \mathbf{r}_j | \mathbf{k} + \mathbf{q}, l' \rangle \langle \mathbf{k} + \mathbf{q}, l' | \mathbf{r}_k | \mathbf{k} + 2\mathbf{q}, l'' \rangle}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}, l} - 2\hbar\omega + 2i\hbar\alpha} \right.$$

$$\times \left. \left(\frac{f(E_{\mathbf{k}+2\mathbf{q}, l''}) - f(E_{\mathbf{k}+\mathbf{q}, l'})}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}+\mathbf{q}, l'} - \hbar\omega + i\hbar\alpha} - \frac{f(E_{\mathbf{k}+\mathbf{q}, l'}) - f(E_{\mathbf{k}, l})}{E_{\mathbf{k}+\mathbf{q}, l'} - E_{\mathbf{k}, l} - \hbar\omega + i\hbar\alpha} \right) \right\}$$

$$\times \left[1 + 4\pi e^2 \sum_{ab} m_a m_b \sum_{\mathbf{k}, l, l''} \langle \mathbf{k}, l | \mathbf{r}_a | \mathbf{k} + 2\mathbf{q}, l'' \rangle \langle \mathbf{k} + 2\mathbf{q}, l'' | \mathbf{r}_b | \mathbf{k}, l \rangle \frac{f(E_{\mathbf{k}+2\mathbf{q}, l''}) - f(E_{\mathbf{k}, l})}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}, l} - 2\hbar\omega + 2i\hbar\alpha} \right]^{-1}$$

Nonlinear intensity

$$I(2\omega; \theta, \Phi, \varphi, \gamma) = \epsilon_0 c_0 |E(2\omega; \theta, \Phi, \varphi, \gamma)|^2$$

Details of the calculation

- WIEN97 spin-polarized scf cycle
 - energy convergence to 0.1 mRy
 - no inversion symmetry (surface)
- add magnetization direction
- our implementation of the spin-orbit coupling
 - second variation
 - spherical potential, inside MT sphere only
 - WIEN97 **basis set**
 - added after complete scf-cycle
- transition matrix elements $\langle \mathbf{k}, l | \mathbf{r}_i | \mathbf{k}', l' \rangle$
- screened nonlinear susceptibility $\chi_{ijk}^{(2)}$
- optical properties ($E, I, \mathcal{A}, \phi_K$)

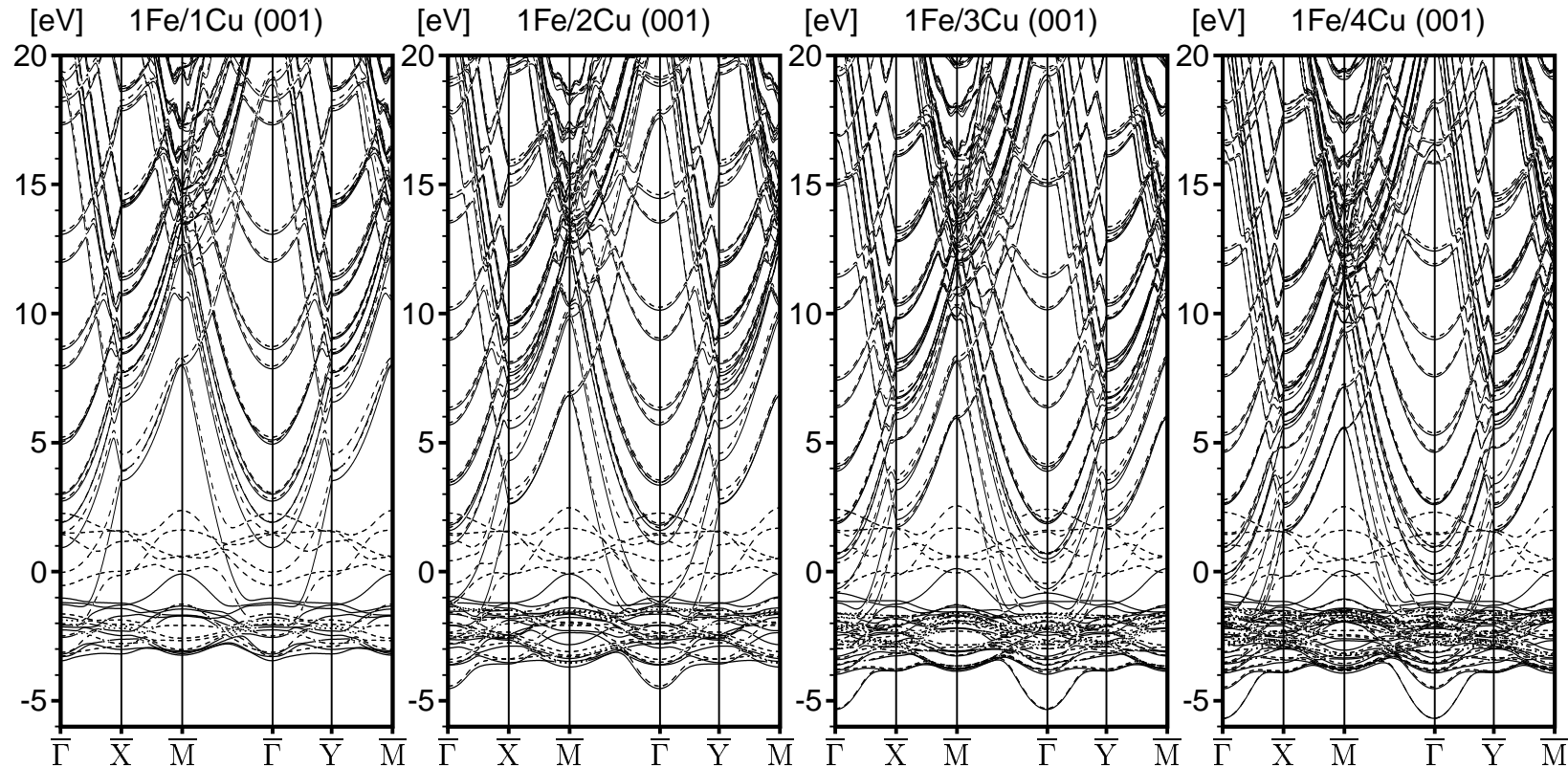
Parameters in WIEN97 scf cycle

- fcc Cu lattice, $a = 3.61 \text{ \AA}$
- separation 8 vacuum layers
- muffin-tin radius 1.21 \AA
- 2502 \mathbf{k} -points in I2BZ

Parameters in optics

- $|E_0(\omega)| = 10^8 \text{ V/m}$
- $\delta z = 4.82 \text{ \AA}$
- $\gamma = 0, \theta = \pi/4$
- linear Cu $n + ik$ from [1]

Fe/Cu (001) Band structure with spin-orbit coupling



Solid lines: > 90% majority spin
Dashed lines: > 90% minority spin
Dotted lines: < 90% of either spin

Magnetic spin moments

1 ML Fe on up to 7 ML fcc Cu (001)

Fe	Cu layer number						
	1	2	3	4	5	6	7
2.828	0.046						
2.844	0.043	-0.026					
2.793	0.033	-0.025	-0.003				
2.792	0.048	-0.015	-0.006	-0.065			
2.853	0.042	-0.016	-0.000	0.001	0.001		
2.835	0.044	-0.016	-0.002	-0.001	-0.002	-0.006	
2.740	0.045	-0.016	0.004	0.005	0.008	0.001	-0.010

FM **AFM**

First Cu layer FM coupled to Fe layer

Second Cu layer AFM coupled to first Cu layer

Nonlinear optical field

$$E(2\omega; \theta, \Phi, \varphi, \gamma) = 2i\delta z \frac{\omega}{c} |E_0(\omega)|^2 \chi_{ijk}^{(2)}(2\mathbf{q}, 2\omega), \mathbf{M} \parallel \mathbf{z}$$

$$\times \left\{ A_s \sin \Phi \begin{pmatrix} -\sin \gamma \\ \cos \gamma \\ 0 \end{pmatrix} + A_p \cos \Phi \begin{pmatrix} F_c \cos \gamma \\ F_c \sin \gamma \\ N^2 F_s \end{pmatrix} \right\} \cdot \begin{pmatrix} 0 & 0 & 0 & | & xyz^- & xxz^+ & 0 \\ 0 & 0 & 0 & | & xxz^+ & -xyz^- & 0 \\ zxx^+ & zxx^+ & zzz^+ & | & 0 & 0 & 0 \end{pmatrix}$$

$$\cdot \begin{pmatrix} [t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma]^2 \\ [t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma]^2 \\ t_p^2 f_s^2 \cos^2 \varphi \\ 2[t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma] t_p f_s \cos \varphi \\ 2[t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma] t_p f_s \cos \varphi \\ 2[t_p f_c \cos \varphi \cos \gamma - t_s \sin \varphi \sin \gamma][t_p f_c \cos \varphi \sin \gamma + t_s \sin \varphi \cos \gamma] \end{pmatrix}$$

Screened nonlinear optical susceptibility

$$\chi_{ijk}^{(2)}(2\mathbf{q}, 2\omega) = e^3 \sum_{\mathbf{k}, l, l''} \left\{ \frac{\langle \mathbf{k} + 2\mathbf{q}, l'' | \mathbf{r}_i | \mathbf{k}, l \rangle \langle \mathbf{k}, l | \mathbf{r}_j | \mathbf{k} + \mathbf{q}, l' \rangle \langle \mathbf{k} + \mathbf{q}, l' | \mathbf{r}_k | \mathbf{k} + 2\mathbf{q}, l'' \rangle}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}, l} - 2\hbar\omega + 2i\hbar\alpha} \right.$$

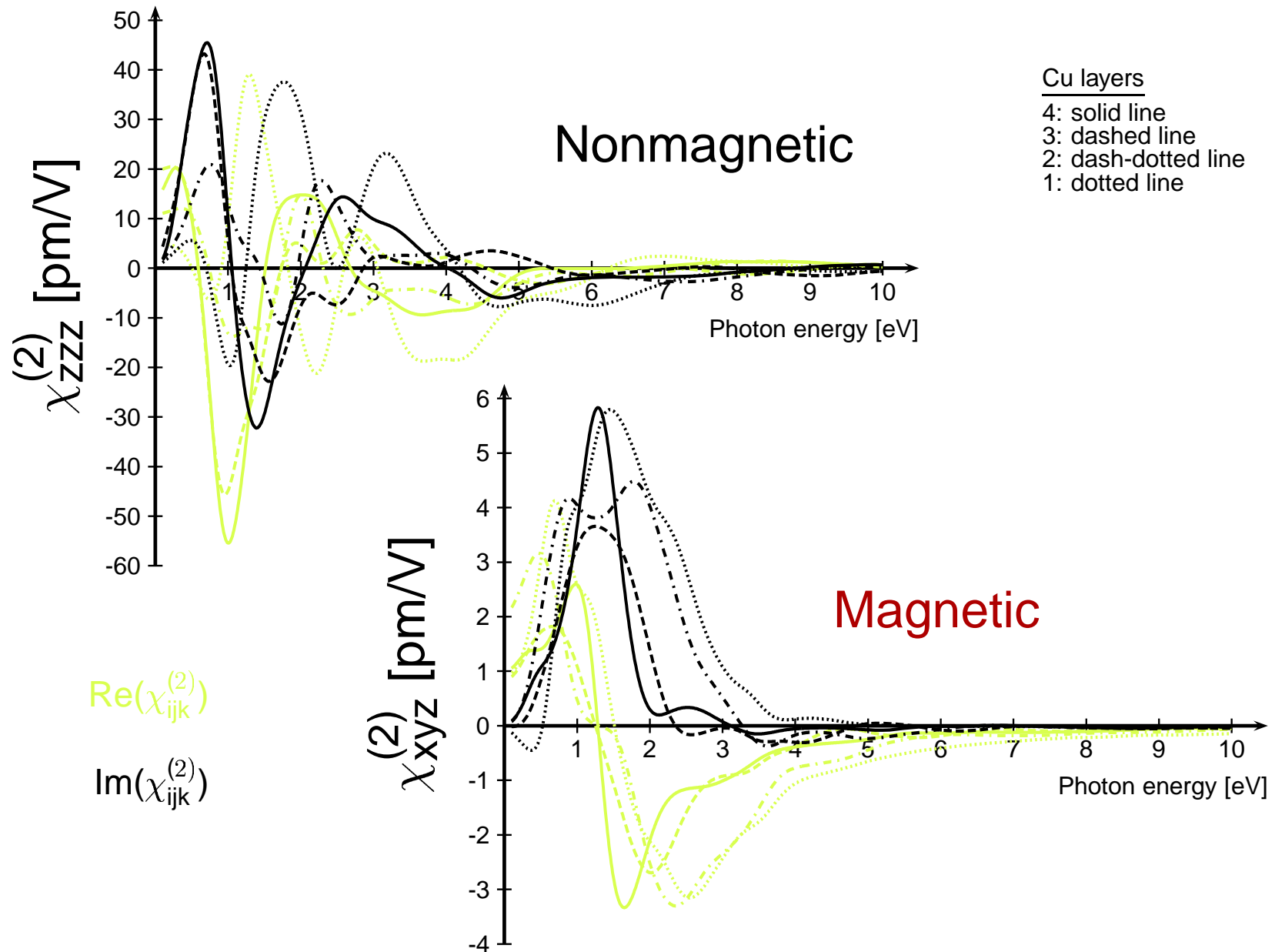
$$\times \left(\frac{f(E_{\mathbf{k}+2\mathbf{q}, l''}) - f(E_{\mathbf{k}+\mathbf{q}, l'})}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}+\mathbf{q}, l'} - \hbar\omega + i\hbar\alpha} - \frac{f(E_{\mathbf{k}+\mathbf{q}, l'}) - f(E_{\mathbf{k}, l})}{E_{\mathbf{k}+\mathbf{q}, l'} - E_{\mathbf{k}, l} - \hbar\omega + i\hbar\alpha} \right) \left. \right\}$$

$$\times \left[1 + 4\pi e^2 \sum_{ab} m_a m_b \sum_{\mathbf{k}, l, l''} \langle \mathbf{k}, l | \mathbf{r}_a | \mathbf{k} + 2\mathbf{q}, l'' \rangle \langle \mathbf{k} + 2\mathbf{q}, l'' | \mathbf{r}_b | \mathbf{k}, l \rangle \frac{f(E_{\mathbf{k}+2\mathbf{q}, l''}) - f(E_{\mathbf{k}, l})}{E_{\mathbf{k}+2\mathbf{q}, l''} - E_{\mathbf{k}, l} - 2\hbar\omega + 2i\hbar\alpha} \right]^{-1}$$

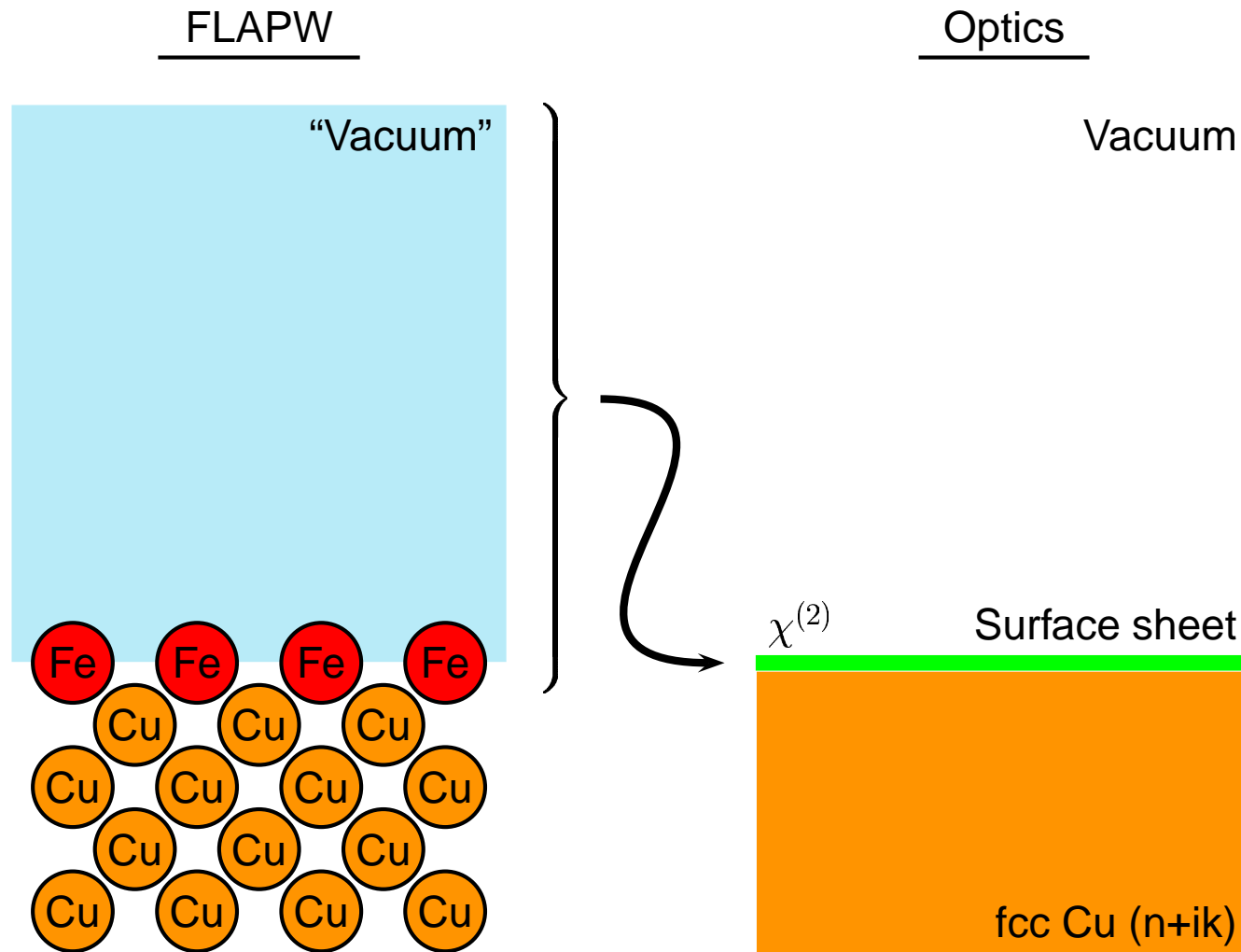
Nonlinear intensity

$$I(2\omega; \theta, \Phi, \varphi, \gamma) = \epsilon_0 c_0 |E(2\omega; \theta, \Phi, \varphi, \gamma)|^2$$

Screened nonlinear susceptibility tensor

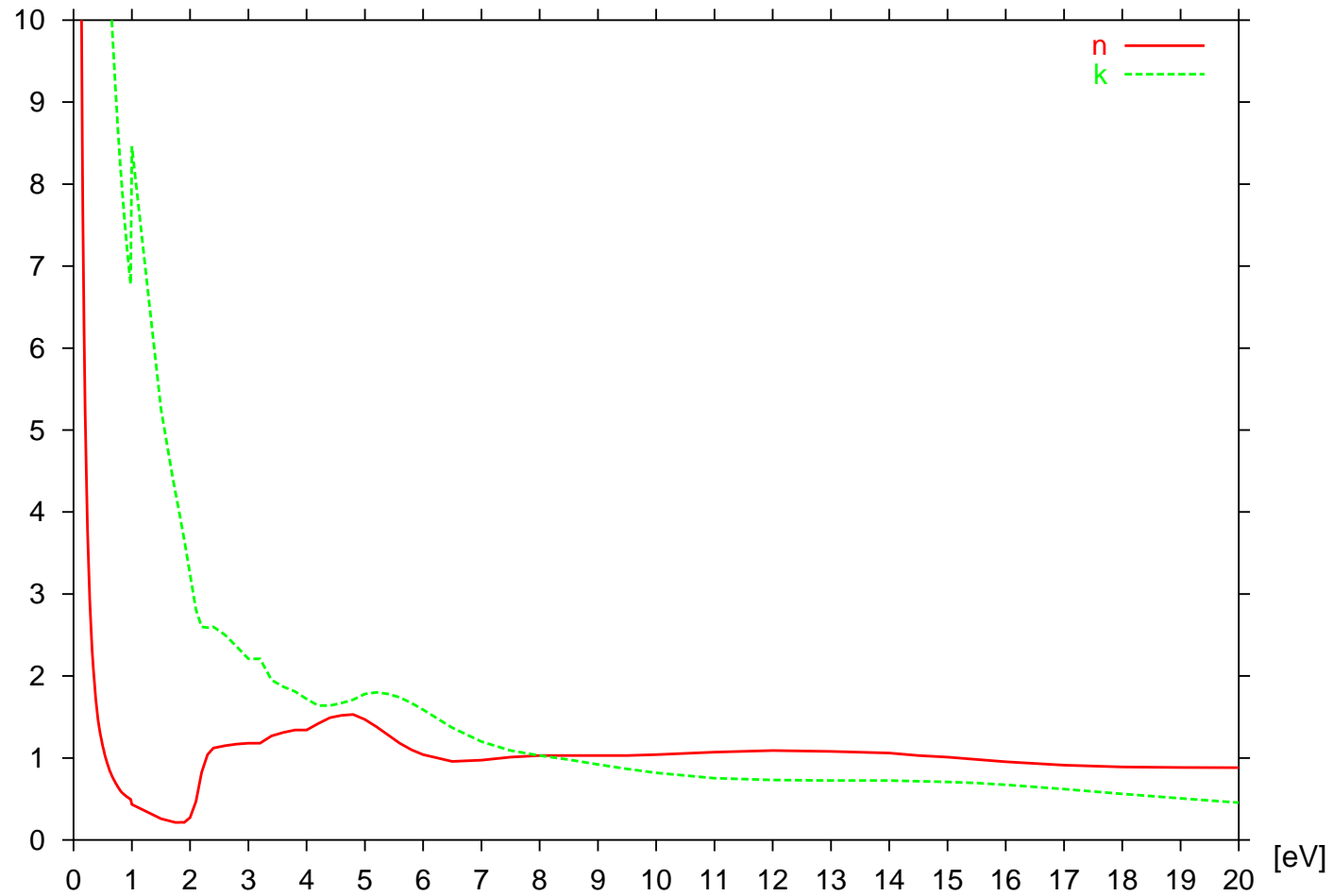


Nonlinear magneto-optics – method



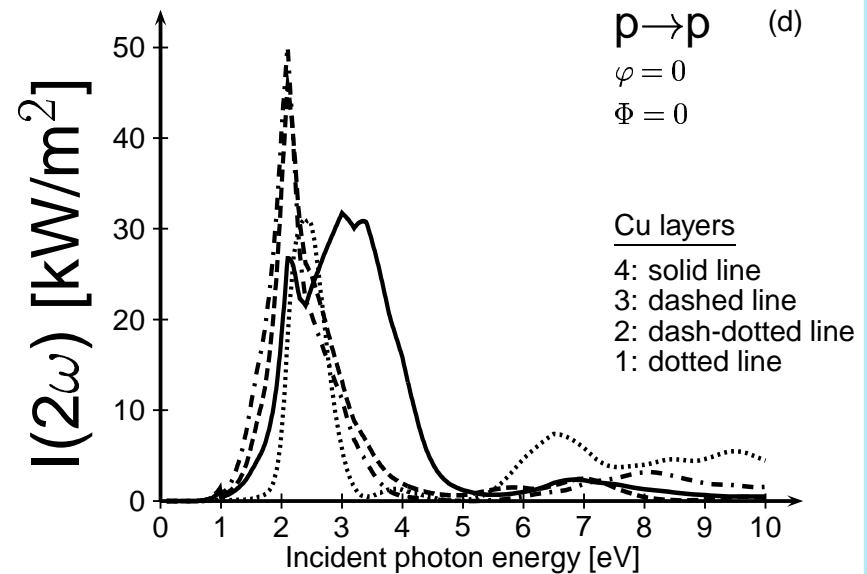
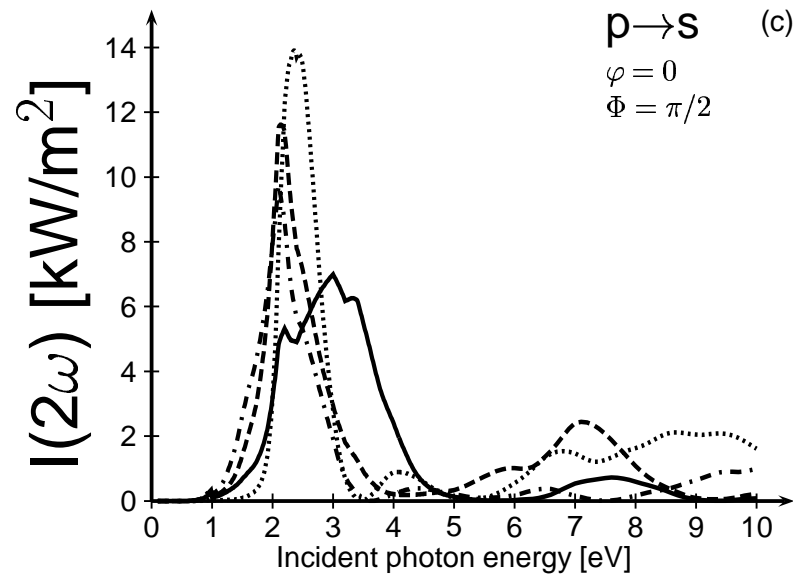
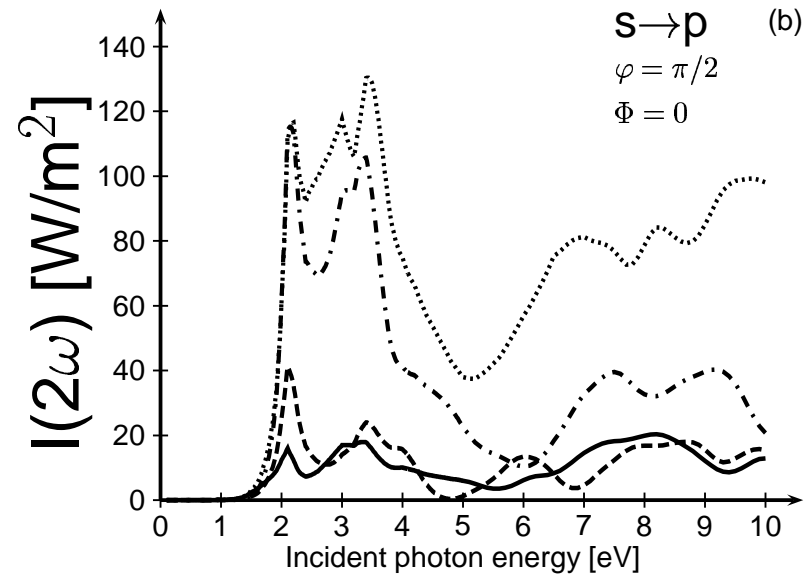
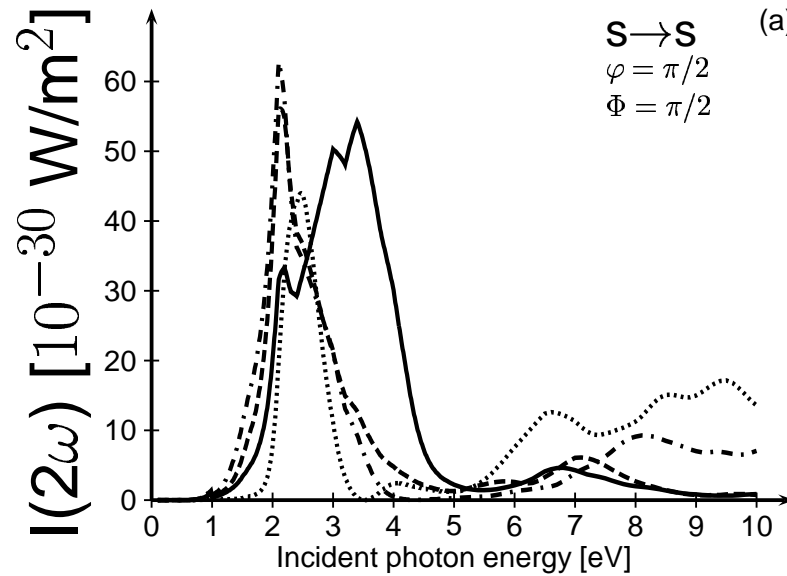
Complex refractive index ($n+ik$) of Cu

Experimental values



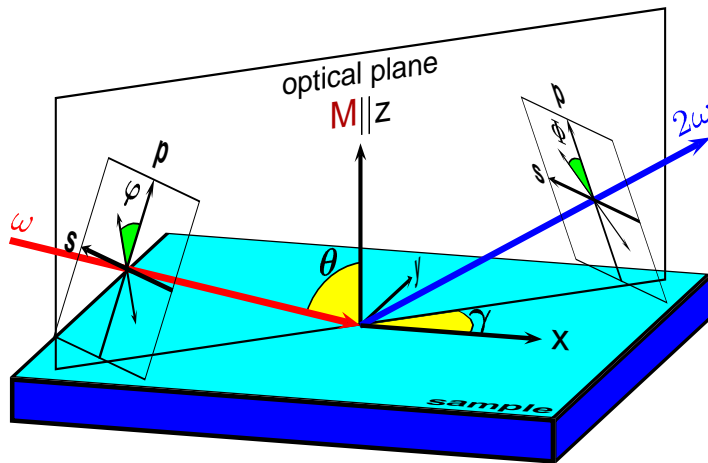
[D. W. Lynch and W. R. Hunter, in *Handbook of optical constants in solids*, edited by E. D. Palik (Academic Press, London, 1985), pp. 275–367]

Second-harmonic intensities



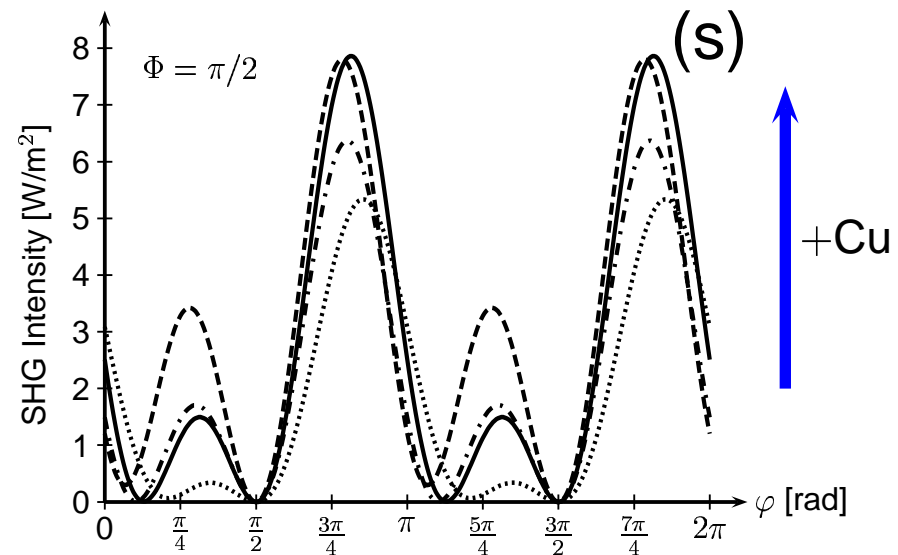
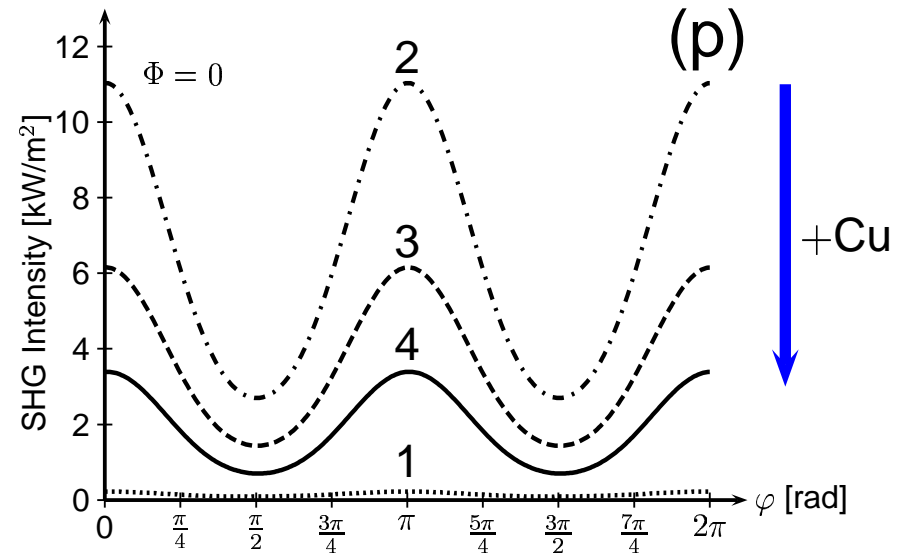
Varying the polarizer angle φ

$$\hbar\omega = 1.5 \text{ eV}$$



Cu layers

- 4: solid line
- 3: dashed line
- 2: dash-dotted line
- 1: dotted line



Nonlinear (2ω) magneto-optical Kerr rotation

$$\phi^{(2)} = \frac{1}{2} \arctan \frac{2\text{Re}[E_s(2\omega)/E_p(2\omega)]}{1 - |E_s(2\omega)/E_p(2\omega)|^2} + \phi_0$$

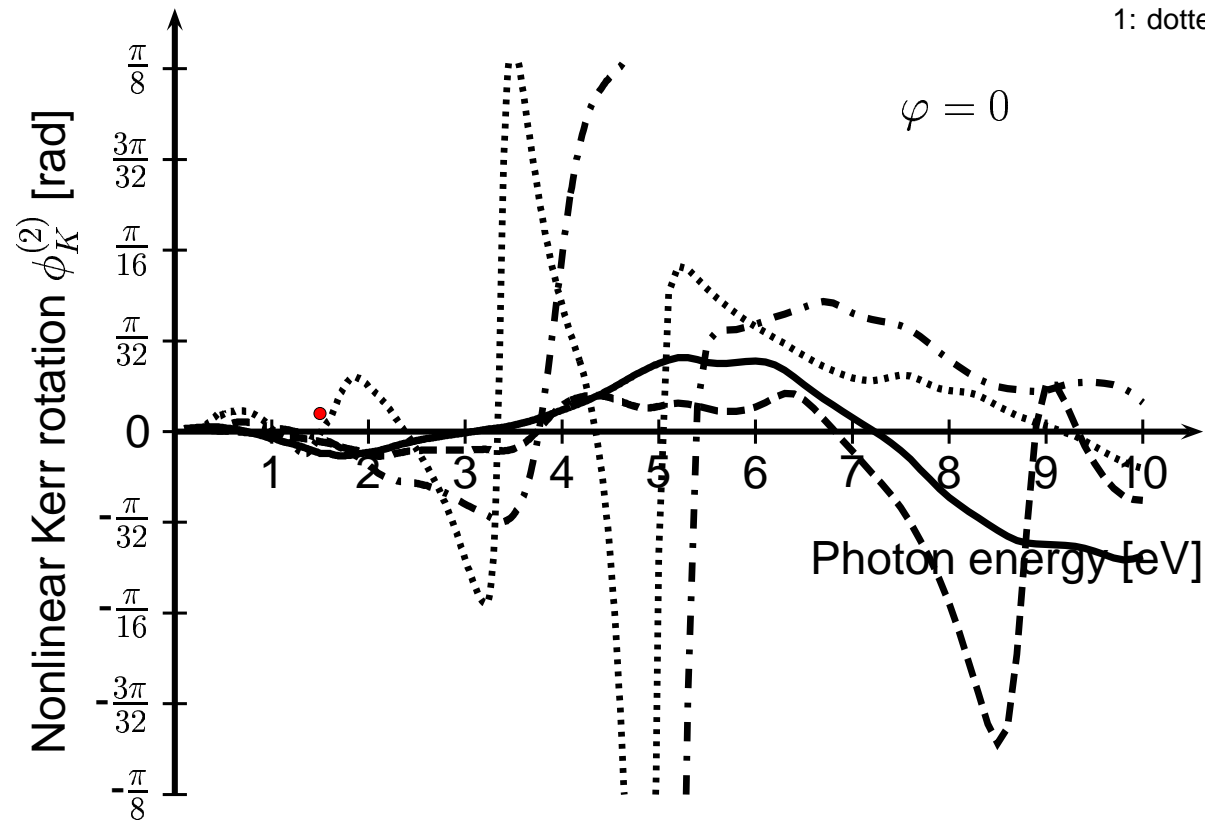
Cu layers

4: solid line

3: dashed line

2: dash-dotted line

1: dotted line



Conclusions

- large magnetic tensor element – powerful tool
- adding Cu gives another degree of freedom
- nonlinear optics can be calculated using first principles methods (WIEN97)
- symmetry breaking by magnetism can be resolved
- susceptibilities for 3 and 4 Cu underlayers close
- intensities do not look that converged for 4 Cu layers
- Kerr rotation not converged

Further details in: Torsten Andersen and W. Hübner, *Substrate effects in magneto-optical second-harmonic generation from first principles: Fe/Cu(001)*, to appear in Phys. Rev. B.

Outlook

- Wien2k spin-orbit?
- correct occupied d-band and excited state positions – GW
- nonspherical contributions to spin-orbit interaction
- nonlocality ($|\mathbf{q}| \neq 0$)
- removal of translational invariance along z

Nonlocal theory of nonlinear magneto-optics

Mesoscopic systems

- Finite size
- Surfaces
- Interfaces
- Short light pulses

Magnetism

- Spin-orbit coupling
- Relativistic electrons

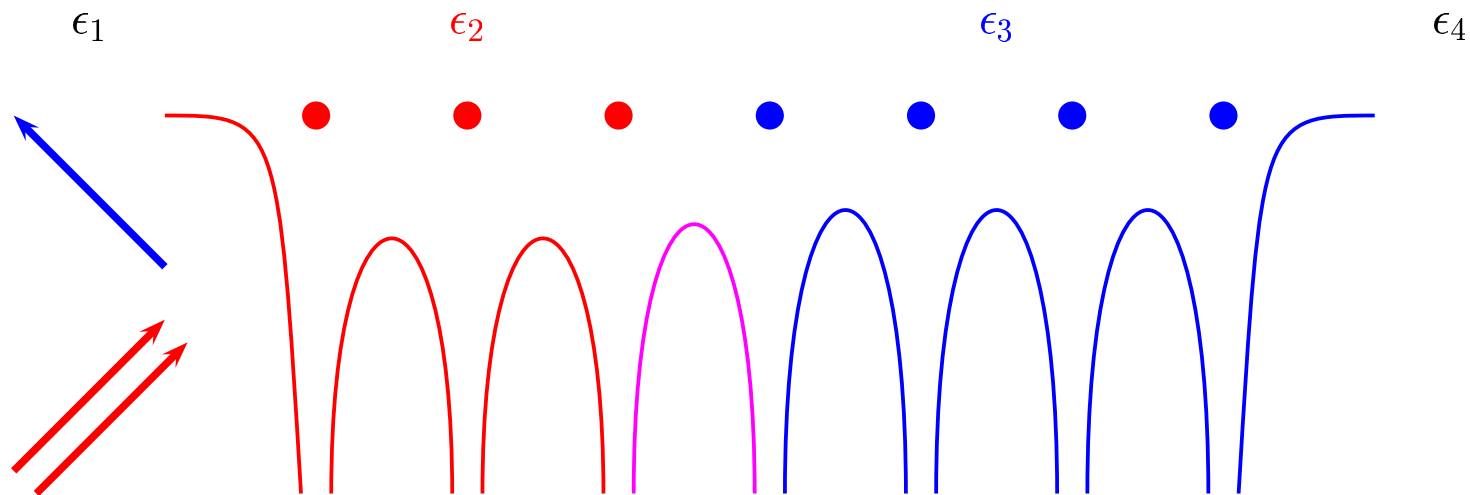
Nonlocality

- Delocalized s -electrons
- Relativistic d -electrons

Problems

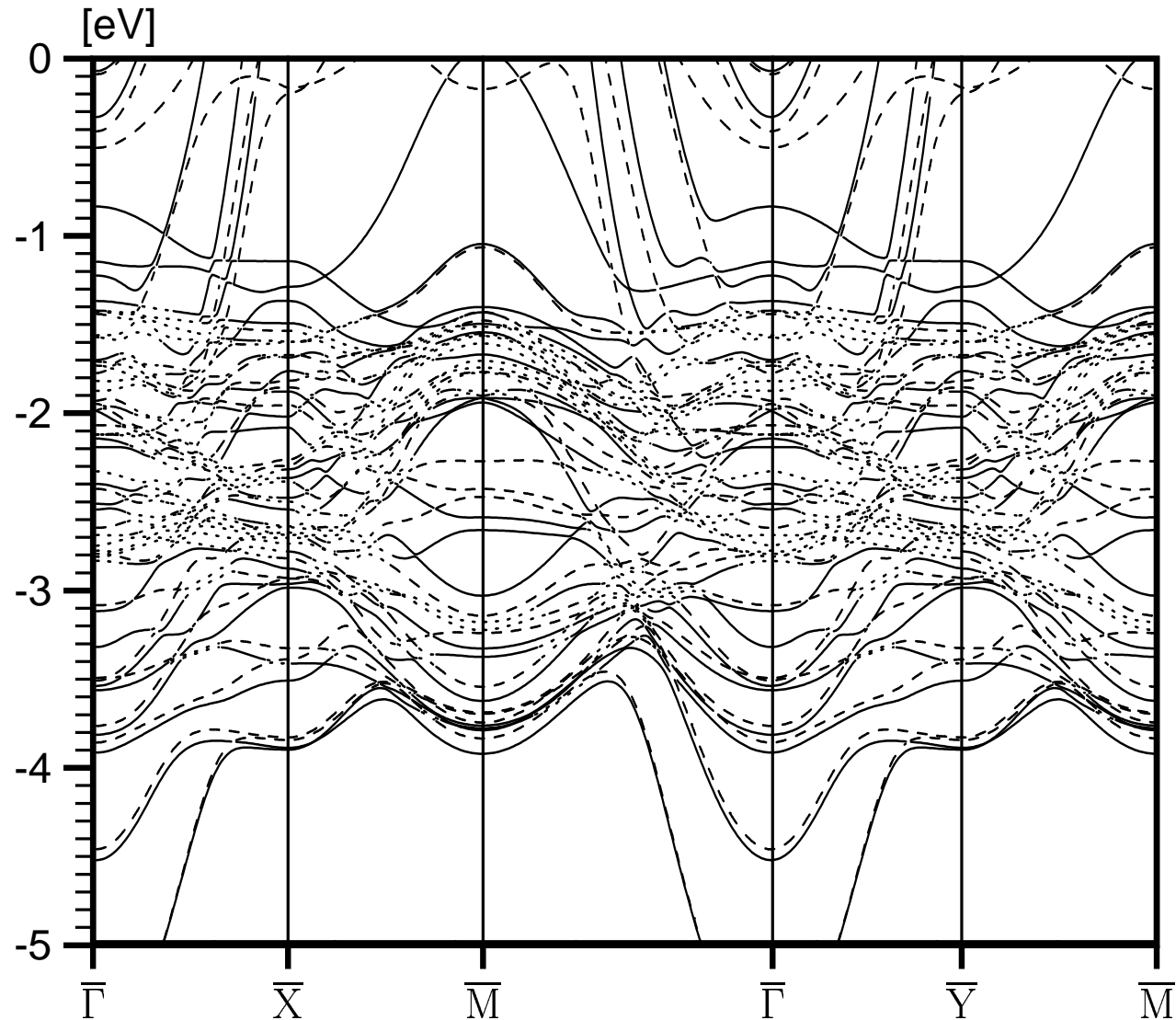
- Boundaries?
- Large field gradients \Rightarrow ~~ED~~
- ED + EQ + MD + ... ?

Nonlinear optics

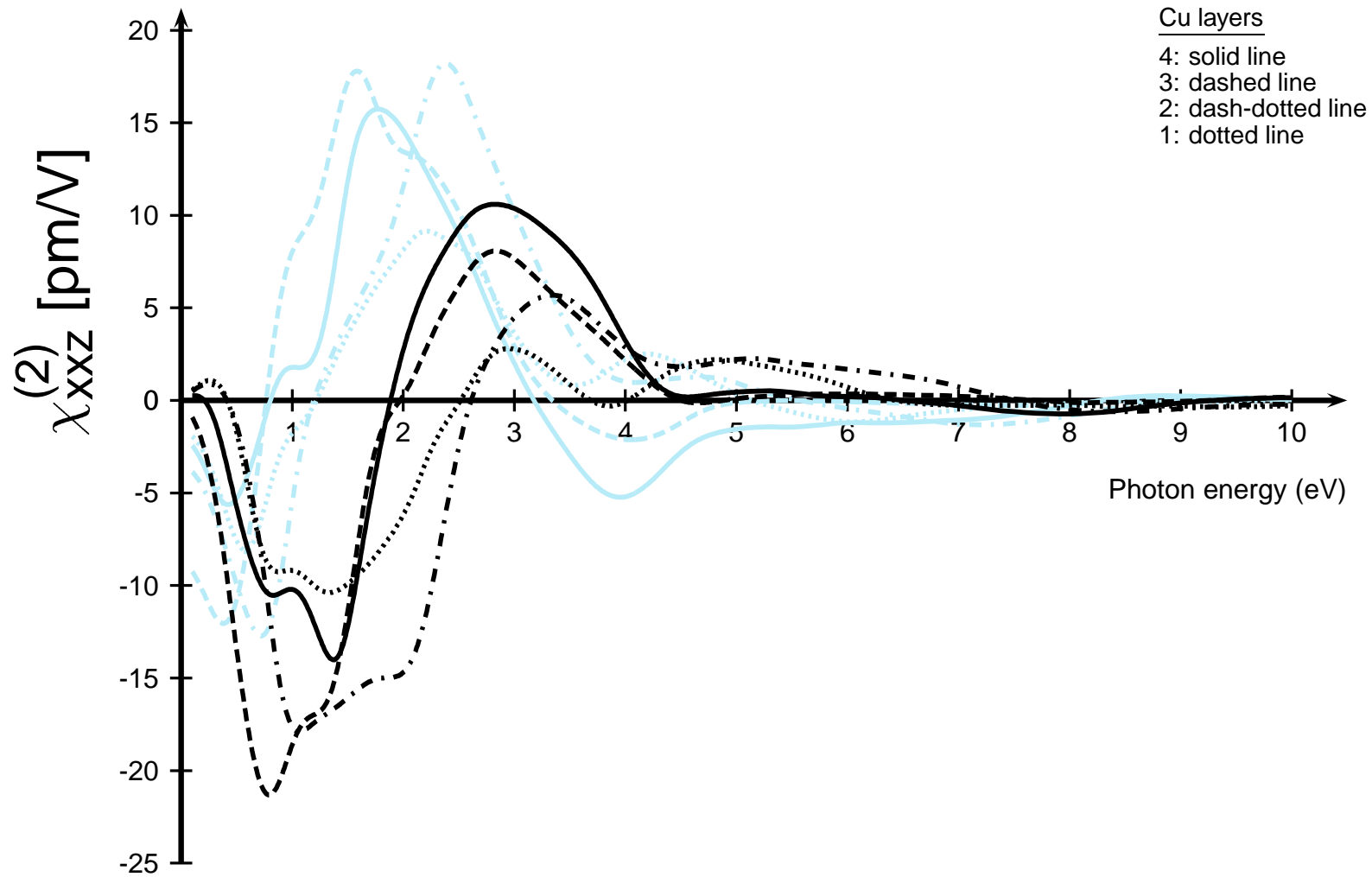


Band structure with spin-orbit coupling

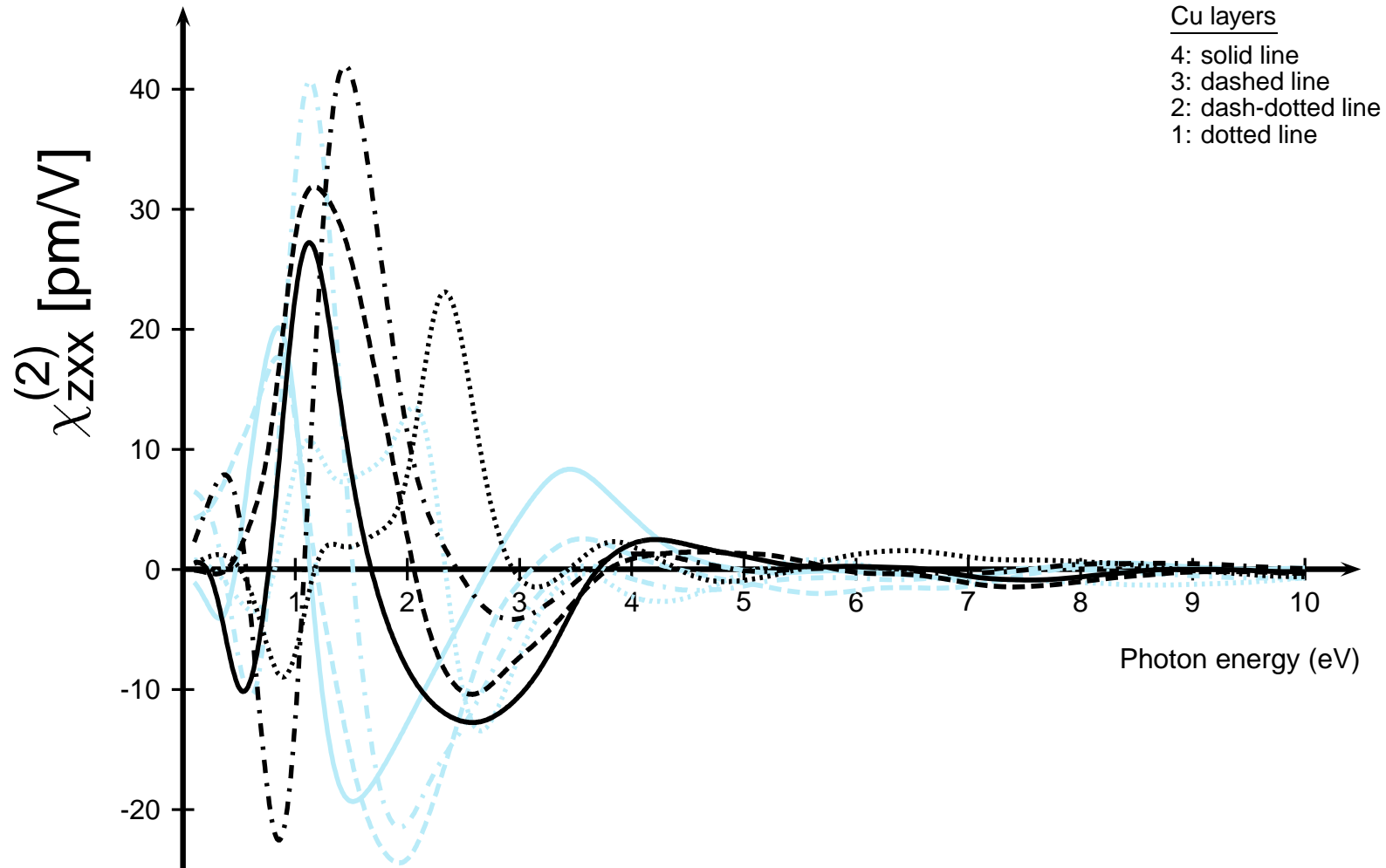
1 ML Fe on 4 ML Cu (001)



Screened nonlinear susceptibility tensor



Screened nonlinear susceptibility tensor



Magneto-optical asymmetry \mathcal{A} as function of φ

$$\mathcal{A} = \frac{I(\mathbf{M} \parallel \mathbf{z}) - I(\mathbf{M} \parallel -\mathbf{z})}{I(\mathbf{M} \parallel \mathbf{z}) + I(\mathbf{M} \parallel -\mathbf{z})}$$

$$\hbar\omega = 1.5 \text{ eV}$$

Cu layers

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